Partitioned multirate domain decomposition waveform relaxation methods for the heat equation

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Outline

Motivation

2 A multirate approach

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A multirate approach

Applications of thermal FSI



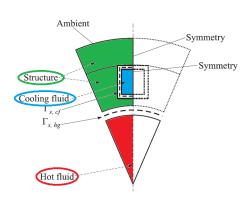
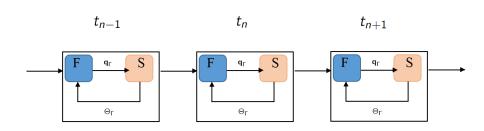


Figure: Left: Vulcain engine for Ariane 5. Right: Sketch of the cooling system; Pline, Wikimedia Commons

Thermal interaction between fluid and structure needs to be modelled

Dirichlet-Neumann coupling

- Partitioned approach for the solution of the coupled problem.
- Fluid Model: Compressible Navier-Stokes FVM.
- Structure Model: Nonlinear heat equation FEM.



Limitations

- The subsolvers are sequential.
- Same time integration for both fields.

Limitations of the Dirichlet-Neumann method

List of wishes

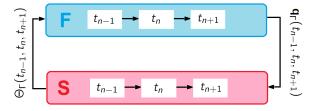
- Two independent time integration schemes.
- High order resolution (at least 2nd order).
- To be able to insert time adaptivity in the framework.

Option 1

Exchange fixed point iteration with time recursion

Option 2

Use a different domain decomposition method



Limitations of the Dirichlet-Neumann method

Option 1

Dirichlet-Neumann Waveform Relaxation (DNWR) algorithm

- + Computationally cheap
- Sequential method

Option 2

Neumann-Neumann Waveform Relaxation (NNWR) algorithm

- + Parallel method
- Computationally expensive

 DNWR and NNWR introduced by Gander and Kwok 2016: Constant coefficients and one single time integration scheme.

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A multirate approach

Model Problem: Coupled heat equations

$$\alpha_{m} \frac{\partial u_{m}(\mathbf{x}, t)}{\partial t} - \nabla \cdot (\lambda_{m} \nabla u_{m}(\mathbf{x}, t)) = 0,$$

$$t \in [t_{0}, t_{f}], \quad \mathbf{x} \in \Omega_{m} \subset \mathbb{R}^{d}, \quad m = 1, 2$$

$$u_{m}(\mathbf{x}, t) = 0, \quad \mathbf{x} \in \partial \Omega_{m} \setminus \Gamma$$

$$u_{1}(\mathbf{x}, t) = u_{2}(\mathbf{x}, t), \quad \mathbf{x} \in \Gamma$$

$$\lambda_{2} \frac{\partial u_{2}(\mathbf{x}, t)}{\partial \mathbf{n}_{2}} = -\lambda_{1} \frac{\partial u_{1}(\mathbf{x}, t)}{\partial \mathbf{n}_{1}}, \quad \mathbf{x} \in \Gamma$$

$$u_{m}(\mathbf{x}, 0) = g_{m}(\mathbf{x}) \quad \mathbf{x} \in \Omega_{m}$$

where $\alpha_m = \lambda_m/D_m$.

Dirichlet-Neumann waveform relaxation (DNWR)

$$\begin{split} & (D) \begin{cases} \alpha_1 \frac{\partial u_1^{k+1}(\mathbf{x},t)}{\partial t} - \nabla \cdot \left(\lambda_1 \nabla u_1^{k+1}(\mathbf{x},t)\right) = 0, & \mathbf{x} \in \Omega_1, \\ u_1^{k+1}(\mathbf{x},t) = 0, & \mathbf{x} \in \partial \Omega_1 \backslash \Gamma, \\ u_1^{k+1}(\mathbf{x},t) = g^k(\mathbf{x},t), & \mathbf{x} \in \Gamma, \\ u_1^{k+1}(\mathbf{x},0) = u_1^0(\mathbf{x}), & \mathbf{x} \in \Omega_1. \end{cases} \\ & (N) \begin{cases} \alpha_2 \frac{\partial u_2^{k+1}(\mathbf{x},t)}{\partial t} - \nabla \cdot \left(\lambda_2 \nabla u_2^{k+1}(\mathbf{x},t)\right) = 0, & \mathbf{x} \in \Omega_2, \\ u_2^{k+1}(\mathbf{x},t) = 0, & \mathbf{x} \in \partial \Omega_2 \backslash \Gamma, \\ \lambda_2 \frac{\partial u_2^{k+1}(\mathbf{x},t)}{\partial \mathbf{n}_2} = -\lambda_1 \frac{\partial u_1^{k+1}(\mathbf{x},t)}{\partial \mathbf{n}_1}, & \mathbf{x} \in \Gamma, \\ u_1^{k+1}(\mathbf{x},0) = u_1^0(\mathbf{x}), & \mathbf{x} \in \Omega_1. \end{cases} \\ & (U) & g^{k+1}(\mathbf{x},t) = \Theta u_2^{k+1}(\mathbf{x},t) + (1-\Theta)g^k(\mathbf{x},t), & \mathbf{x} \in \Gamma. \end{cases}$$

How to choose the relaxation parameter ⊖ properly?

Neumann-Neumann waveform relaxation (NNWR)

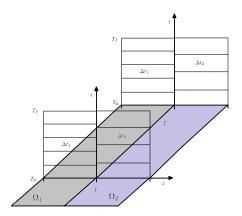
$$(D_{m}), m = 1, 2 \begin{cases} \alpha_{m} \frac{\partial u_{m}^{k+1}(\mathbf{x}, t)}{\partial t} - \nabla \cdot (\lambda_{m} \nabla u_{m}^{k+1}(\mathbf{x}, t)) = 0, & \mathbf{x} \in \Omega_{m}, \\ u_{m}^{k+1}(\mathbf{x}, t) = 0, & \mathbf{x} \in \partial \Omega_{m} \backslash \Gamma, \\ u_{m}^{k+1}(\mathbf{x}, t) = g^{k}(\mathbf{x}, t), & \mathbf{x} \in \Gamma, \\ u_{m}^{k+1}(\mathbf{x}, 0) = u_{1}^{0}(\mathbf{x}), & \mathbf{x} \in \Omega_{m}. \end{cases}$$

$$(N_{m}), m = 1, 2 \begin{cases} \alpha_{m} \frac{\partial \psi_{m}^{k+1}(\mathbf{x}, t)}{\partial t} - \nabla \cdot (\lambda_{m} \nabla \psi_{m}^{k+1}(\mathbf{x}, t)) = 0, & \mathbf{x} \in \Omega_{m}, \\ \psi_{m}^{k+1}(\mathbf{x}, t) = 0, & \mathbf{x} \in \partial \Omega_{m} \backslash \Gamma, \\ \lambda_{m} \frac{\partial \psi_{m}^{k+1}(\mathbf{x}, t)}{\partial \mathbf{n}_{1}} = \lambda_{1} \frac{\partial u_{1}^{k+1}(\mathbf{x}, t)}{\partial \mathbf{n}_{1}} + \lambda_{2} \frac{\partial u_{2}^{k+1}(\mathbf{x}, t)}{\partial \mathbf{n}_{2}}, & \mathbf{x} \in \Gamma, \\ \psi_{m}^{k+1}(\mathbf{x}, 0) = 0, & \mathbf{x} \in \Omega_{m}. \end{cases}$$

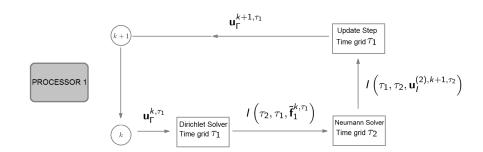
How to choose the relaxation parameter ⊖ properly?

Choices

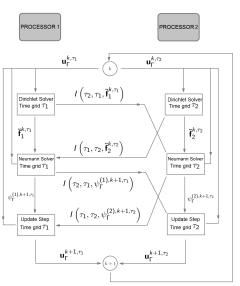
- Space discretization: 1D and 2D finite elements
- Time discretization: Implicit Euler and SDIRK2
- Matching space grid at the interface, unknowns on interface
- Nonmatching time grids at the interface, linear interpolation



Multirate DNWR Algorithm



Multirate NNWR Algorithm



Relaxation parameter

- Space discretization: 1D equidistant FE/FE.
- Time integration: nonmultirate Implicit Euler.
- Use eigendecomposition of the tridiagonal Toeplitz matrices $\mathbf{M}/\Delta t + \mathbf{A}$ to compute the spectral radius of the iteration matrix w.r.t. \mathbf{u}_{Γ} .

DNWR Algorithm

NNWR Algorithm

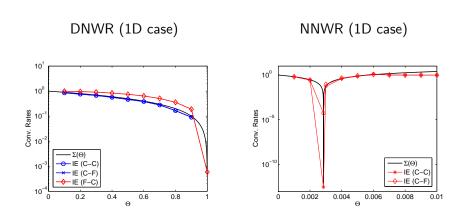
$$\Theta_{opt} = \left(1 + \frac{\mathbf{S}^{(1)}}{\mathbf{S}^{(2)}}\right)^{-1}, \qquad \qquad \Theta_{opt} = \left(2 + \frac{\mathbf{S}^{(1)}}{\mathbf{S}^{(2)}} + \frac{\mathbf{S}^{(2)}}{\mathbf{S}^{(1)}}\right)^{-1},$$

with

$$\mathbf{S}^{(m)} = (6\Delta x(\alpha_m \Delta x^2 + 3\lambda_m \Delta t) - (\alpha_m \Delta x^2 - 6\lambda_m \Delta t)^2 s_m).$$

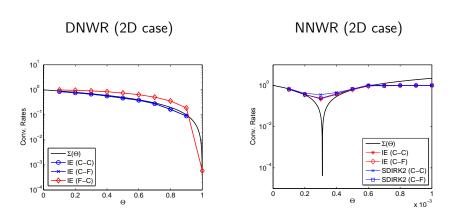
 Θ_{opt} gives the optimal parameter for any coupled materials!

Convergence Rates: Air-Water coupling



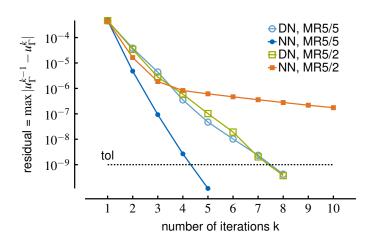
 $\Theta_{\textit{opt}}$ hits the optimal convergent performance.

Convergence Rates: Air-Water coupling



The NNWR method is more sensitive than the DNWR method.

Comparison DNWR and NNWR



NNWR uses half the iterations than DNWR but is less robust!

Conclusions

- Multirate parallel method for coupled parabolic problems (NNWR).
- Multirate sequential method for coupled parabolic problems (DNWR).
- We have performed a 1D analysis to find Θ_{opt} .
- Θ_{opt} is dependent on Δt , Δx , λ_m and α_m , m=1,2.
- Θ_{opt} works for 2D, multirate and time adaptivity.
- Θ_{opt} is more sensitive for NNWR than for DNWR.

More at:

- A. Monge, P. Birken, SISC, accepted (arXiv:1805.04336)
- A. Monge, P. Birken, Proceedings of 25th Domain Decomposition Conference, accepted.
- A. Monge, P. Birken, PAMM 19, submitted.



Thank you!















